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“Acceleration and Radiation”

6-week progress report

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Outline

- * Is there a problem?
- * Lorentz-Abraham-Dirac (LAD) Equation
- * Runaway & pre-acceleration solutions
- * Other approaches
- * Example: Model dependent physics predictions in oscillating electric field
- * Towards a consistent picture: Radiation in Kaluza theory
- * Future objectives

The Inconsistency in Lorentz Force

The Lorentz force $m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta$

Larmor showed that the radiation rate of an accelerated charge is

$$\frac{dE_R}{dt} = \frac{2}{3}e^2\dot{v}^2$$

Relativistically (Heaviside-Abraham) the lost of energy-momentum due to radiation is

$$\dot{P}^\alpha = \frac{2}{3}e^2 (\dot{u}^\beta \dot{u}_\beta) u^\alpha$$

Or using $0 = \frac{d}{d\tau} (\dot{u}^\beta u_\beta) = \ddot{u}^\beta u_\beta + \dot{u}^\beta \dot{u}_\beta$ we get

$$\dot{P}^\alpha = -\frac{2}{3}e^2 (\ddot{u}^\beta u_\beta) u^\alpha$$

An accelerating charge loses energy! An effect that is missing in Lorentz force.

Lorentz-Abraham-Dirac equation

Derivation by Dirac

Dirac uses the retarded Liénard-Wiechert potential (solution of Maxwell equations), conservation of energy and “simplicity” to derive the covariant Lorentz-Abraham equation

$$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + m\tau_0 (\delta_\beta^\alpha + u^\alpha u_\beta) \ddot{u}^\beta$$

“But whereas these equations, as derived from the Lorentz theory, are only approximate, we now see that there is good reason for believing them to be exact, within the classical theory.”

Dirac assumes the simplest possible form of energy flow, and justifies as follows:

“[Other forms] are all much more complicated than [the form used], so that one would hardly expect them to apply to a simple thing like an electron.” Dirac, 1938

Runaway solutions

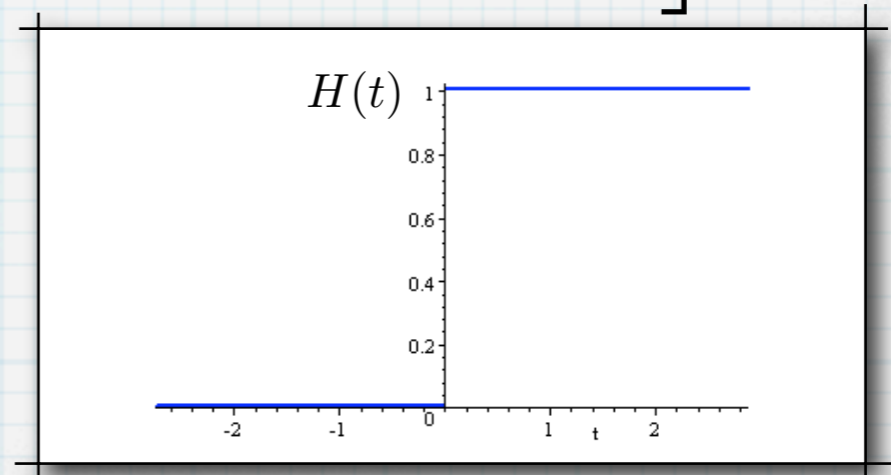
The non-relativistic LAD (F_L is the Lorentz force)

$$a - \tau_0 \dot{a} = \frac{1}{m} F_L$$

has a general solution

$$a(t) = e^{t/\tau_0} \left[a_0 - \frac{1}{m\tau_0} \int_0^t F_L(t') e^{-t'/\tau_0} dt' \right]$$

for an electric field $H(t)E$ that is turned abruptly at $t=0$



$$a(t) = e^{t/\tau_0} \left[a_0 - \frac{eE}{m} \left(1 - e^{-t/\tau_0} \right) H(t) \right]$$

the acceleration grows exponentially!

Pre-acceleration solutions

We exclude runaway solutions by setting

$$a_0 = \frac{eE}{m} [H(t) + H(-t)]$$

and obtain

$$a(t) = \frac{eE}{m} e^{t/\tau_0} H(-t) + \frac{eE}{m} H(t)$$

the particle moves before time $t=0$.

So the effect precedes the cause...

The principle of causality breaks down!

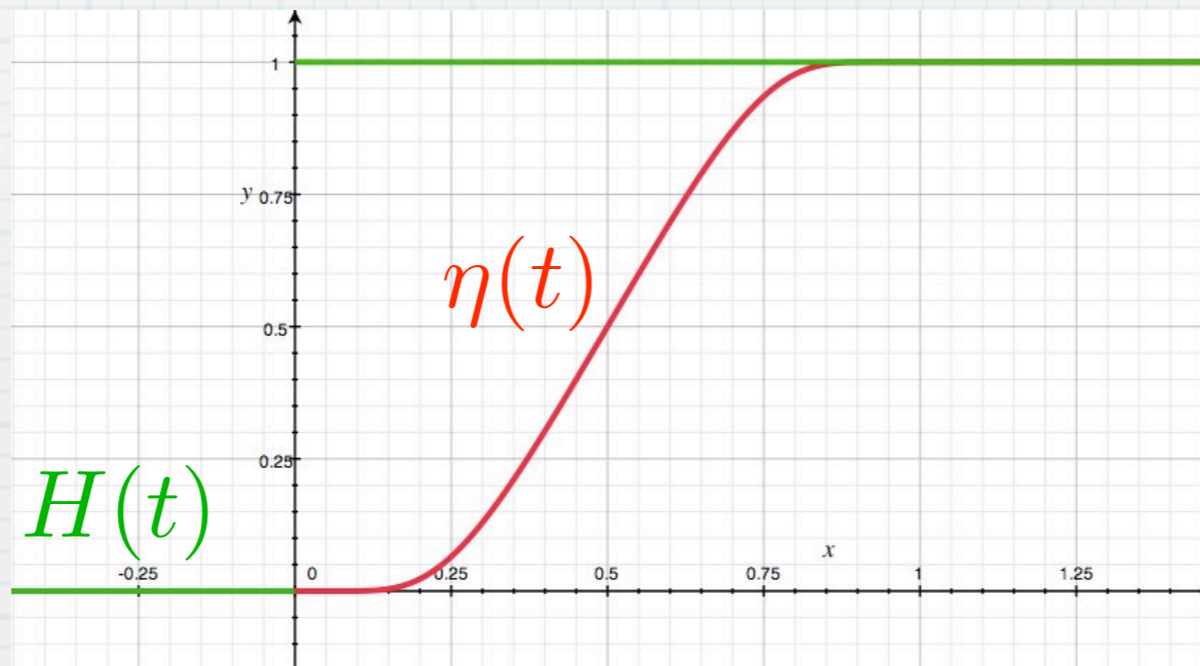
...does it?!

Resolving the Inconsistencies

The derivation of the Lorentz-Abraham-Dirac equation is based on a Taylor expansion of the field.

But a step function $H(t)E$ is non-analytic and cannot be expanded!

Solution: Take a smooth electric field $\eta(t)E$ that vanishes at $t=0$ and assume that the acceleration ceases as $t \rightarrow \infty$



This removes all pathologies (Yaghjian 1992).

Landau-Lifshitz equation

The LAD equation:

$$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + m\tau_0 (\delta_\beta^\alpha + u^\alpha u_\beta) \ddot{u}^\beta$$

If τ_c is a characteristic time scale, the last term is of order τ_0/τ_c

In the first order we get the Lorentz equation:

$$m\dot{u}^\alpha \cong eF^{\alpha\beta}u_\beta$$

Hence,

$$m\ddot{u}^\alpha \cong e \frac{d}{d\tau} (F^{\alpha\beta}u_\beta)$$

Using this approximation in the LAD equation, gives the Landau-Lifshitz equation

$$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + e\tau_0 [F_{,\gamma}^{\alpha\beta}u_\beta u^\gamma + e/m (\delta_\beta^\alpha + u^\alpha u_\beta) F_\gamma^\beta F_\delta^\gamma u^\delta]$$

Mo-Papas equation

The Lorentz equation:

$$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta$$

The radiation-reaction force is (Larmor-Heaviside-Abraham)

$$F_{RR}^\alpha = -\frac{2}{3}e^2\dot{u}_\beta\dot{u}^\beta u^\alpha$$

In the first order approximation

$$F_{RR}^\alpha \cong -\frac{2}{3}\frac{e^3}{m}F^{\beta\gamma}\dot{u}_\beta u_\gamma u^\alpha$$

Postulate: The particle experiences another force $e\tau_0 F^{\alpha\beta}\dot{u}_\beta$

$$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + e\tau_0 (F^{\alpha\beta}\dot{u}_\beta + F^{\beta\gamma}\dot{u}_\beta u_\gamma u^\alpha)$$

Caldirola's equation

Iterations of the LAD give infinitely many derivatives of the velocity, hence a non-local theory.

The non-locality means that the four-velocity $u^\alpha(\tau)$ depends on the four-velocity at a previous time $u^\alpha(\tau - \Delta\tau)$

Postulate a universal quanta of time - the chronon τ_0

Want: A finite-difference relativistic equation for the four-velocity, that in the limit $\tau_0 \rightarrow 0$ gives the LAD

Get Caldirola's equation

$$\frac{m}{\tau_0} \left\{ \cancel{u^\alpha(\tau)} - u^\alpha(\tau - \tau_0) + \underbrace{u^\alpha(\tau) u_\beta(\tau) [\cancel{u^\beta(\tau)} - u^\beta(\tau - \tau_0)]}_{-u^\alpha(\tau)} \right\} = eF^{\alpha\beta}(\tau) u_\beta(\tau)$$

Caldirola-Yaghjian equation

Postulate: the electron has the internal structure of a charged sphere

Compute the self-force due to its internal structure in the particle's rest frame

Lorentz transforming to a general frame, gives

$$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + \underbrace{\frac{m}{\tau_0} \left[u^\alpha(\tau - \tau_0) + u^\alpha(\tau)u^\beta(\tau)u_\beta(\tau - \tau_0) \right]}_{\text{Caldirola term}}$$

Caldirola term

The candidates

LAD (1904-1938)	$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + m\tau_0 (\delta_\beta^\alpha + u^\alpha u_\beta) \ddot{u}^\beta$
Landau-Lifshitz (1951)	$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + e\tau_0 [F_{,\gamma}^{\alpha\beta}u_\beta u^\gamma + e/m (\delta_\beta^\alpha + u^\alpha u_\beta) F_\gamma^\beta F_\delta^\gamma u^\delta]$
Mo-Papas (1971)	$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + e\tau_0 (F^{\alpha\beta}\dot{u}_\beta + F^{\beta\gamma}\dot{u}_\beta u_\gamma u^\alpha)$
Caldirola (1979)	$-m/\tau_0 [u^\alpha (\tau - \tau_0) + u^\alpha (\tau) u_\beta (\tau) u^\beta (\tau - \tau_0)] = eF^{\alpha\beta} (\tau) u_\beta (\tau)$
Caldirola-Yaghjian (1992)	$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + m/\tau_0 [u^\alpha (\tau - \tau_0) + u^\alpha (\tau) u^\beta (\tau) u_\beta (\tau - \tau_0)]$

Example: different physics in an oscillating electric field

Take $\vec{E} = E \cos(\omega t) \hat{x}$ (v denotes 3-velocity)

Lorentz and Mo-Papas eq.	$\dot{v} = \frac{eE}{m} \cos(\omega t)$
LAD	$\dot{v} = \frac{eE}{m} \cos(\omega t) - \frac{2e^2}{3m} \frac{\dot{v}^2 v - \ddot{v}(1+v^2)}{(1+v^2)^{3/2}}$
Landau-Lifshitz	$\dot{v} = \frac{eE}{m} \cos(\omega t) + \frac{eE}{m} \tau_0 \omega \sin(\omega t) \sqrt{1+v^2}$
Caldirola	$v \sqrt{1+v_-^2} - v_- \sqrt{1+v^2} = \frac{eE}{m} \tau_0 \cos(\omega t)$
Caldirola-Yaghjian	$\dot{v} = \frac{eE}{m} \cos(\omega t) - \frac{1}{\tau_0} \left[v \sqrt{1+v_-^2} - v_- \sqrt{1+v^2} \right]$

Kaluza Theory

Consider a five-dimensional space-time, with the metric

$$\hat{g}_{AB} = \begin{bmatrix} g_{ab} + k^2 \phi^2 A_a A_b & k\phi^2 A_a \\ k\phi^2 A_b & \phi^2 \end{bmatrix}$$

Varying the 5D Hilbert-Einstein action

$$S = -\frac{1}{16\pi\hat{G}} \int \hat{R} \sqrt{-\hat{g}} d^5 x$$

Get a unification of Einstein's equation and Maxwell's equations:

$$\begin{aligned} G_{ab} &= \frac{k^2 \phi^2}{2} T_{ab} - \frac{1}{\phi} (\nabla_a (\partial_b \phi) - g_{ab} \square \phi) \\ \nabla^a F_{ab} &= -3 \frac{\partial^a \phi}{\phi} F_{ab} \\ \square \phi &= \frac{k^2 \phi^3}{4} F_{ab} F^{ab} \end{aligned}$$

Back to radiation

By the equivalence principle, gravity and acceleration are intimately related.

Both gravity & electromagnetism enter the radiation effects. That's where the unification scheme of Kaluza comes into play.

The 4D geodesic equation from 5D Kaluza theory is

$$\underbrace{\dot{u}^a + \Gamma_{bc}^a u^b u^c}_{\text{Einstein term}} = \underbrace{\frac{e}{m} F^{ab} u_b}_{\text{Lorentz}} + \underbrace{G_{bc}^a u^b u^c}_{\text{Lifshitz}} + \underbrace{H_b^a u^b u^5 + I^a (u^5)^2}_{\text{Additional terms}}$$

where G , H and I are tensors that depend non-linearly on the fields.

The equation is very similar to the Landau-Lifshitz equation, has no pathological solutions but requires further exploration of the motion in the fifth dimension.

Future objectives

- * Search for generalized action principle leading to existing models
- * Search for an action principle unifying gravity and electromagnetism (example: Kaluza theory)
- * Obtain τ_0^2 terms:

$$m\dot{u}^\alpha = eF^{\alpha\beta}u_\beta + \tau_0 X^\alpha + \tau_0^2 Y^\alpha + \dots$$

References

- * Dirac. Classical theory of radiating electrons. Proceedings of the Royal Society of London. Series A (1938)
- * L. D. Landau and E. M. Lifshitz. The Classical theory of Fields. Pergamon, Oxford (1962)
- * Mo and Papas. New equation of motion for classical charged particles. Physical Review D (1971)
- * Caldirola. A relativistic theory of the classical electron. La Rivista del Nuovo Cimento (1978-1999) (1979)
- * A. D. Yaghjian, Relativistic Dynamics of a Charged Sphere, Lecture Notes in Physics m1 1, Springer-Verlag, Berlin (1992)
- * Rohrlich. The dynamics of a charged sphere and the electron. American Journal of Physics (1997)
- * Poisson. An introduction to the Lorentz-Dirac equation. Arxiv preprint gr-qc (1999)
- * Rohrlich. Dynamics of a charged particle. Physical Review E (2008)
- * Overduin and Wesson. Kaluza-Klein Gravity. arXiv (1998) vol. gr-qc